

# Suppression of Chaos Propagation in Ladder Chaotic Circuits by Local Switching of Coupling Strength

Naoto Yonemoto, Yoko Uwate and Yoshifumi Nishio  
 Dept. of Electrical and Electronic Engineering  
 Tokushima University  
 2-1 Minami-Josanjima, Tokushima 7708506, Japan  
 E-mail: {yonex34, uwate, nishio}@ee.tokushima-u.ac.jp

**Abstract**— Various phenomena have been observed in coupled chaotic circuits. One of the phenomena in chaotic synchronization is the chaotic propagation. It has been investigated what kind of frustrations affect the chaotic propagation. These have been investigated for both simple and complex networks. In this study, the chaos propagation in the ladder-coupled chaotic circuits is studied by computer simulation, using the voltage difference between the adjacent circuits as the threshold for switching the coupling strength.

**Keywords;** Chaotic circuit; Chaos propagation

## I. INTRODUCTION

Various phenomena have been observed in coupled chaotic circuits. The collapse of the chaotic synchronization [1] and the clustering phenomena [2] have attracted attention and have been studied. However, it is believed that there are still many unexplained phenomena in chaotic synchronization. Therefore, it is important to discover, model, and investigate such phenomena in order to understand and utilize them.

One of the synchronous phenomena observed in coupled chaotic circuits is chaos propagation, which occurs when one chaotic circuit is set to generate chaos and the others are set to generate a three-period attractor. The propagation time and the propagation rate have been considered by the coupling strength [3], changing the network structure, and the location where the chaos is generated [4]. In the past, a method to change the network topology by probability [5] has been proposed as one of the frustrations to the chaos synchronization.

In this study, we utilize this method to investigate the effect of switching the coupling strength on the propagation of chaos in ladder chaotic circuits. The chaos propagation is investigated in three different ways of changing the coupling strength by means of computer simulation.

## II. SYSTEM MODEL

The chaotic circuit is shown in Fig. 1. This circuit consists of a negative resistor, two inductors, a capacitor and dual-directional diodes. The system model is shown in Fig. 2. In this system, the circuit in the one end of the system generates chaotic attractor and the other circuits generate three-periodic attractors. Ten chaotic circuits next to each other are coupled together with resistors.

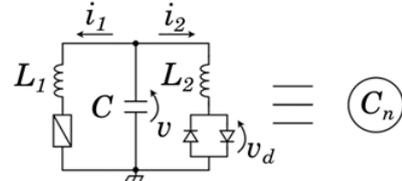


Figure 1: Circuit model.

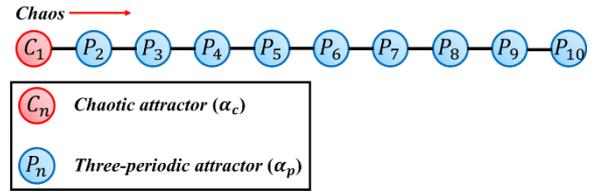


Figure 2: System model.

The parameters are described as follows:

$$\begin{cases} i_1 = \sqrt{\frac{C}{L_1}} V x_n, & i_2 = \frac{\sqrt{L_1 C}}{L_2} V y_n, & v = V z_n \\ \alpha = r \sqrt{\frac{C}{L_1}}, & \beta = \frac{L_1}{L_2}, & \delta = r_d \frac{\sqrt{L_1 C}}{L_2}, \\ \gamma = \frac{1}{R}, & t = \sqrt{L_1 C} \tau. \end{cases}$$

The normalized circuit equations are described as follows:

$$\begin{cases} \frac{dx_n}{d\tau} = \alpha x_n + z_n \\ \frac{dy_n}{d\tau} = z_n - f(y_n) \\ \frac{dz_n}{d\tau} = -x_n - \beta y_n - \gamma_{n,n-1} (z_n - z_{n-1}) \\ \quad - \gamma_{n,n+1} (z_n - z_{n+1}) \\ (n = 1, 2, \dots, N) \end{cases} \quad (1)$$

where

$$\gamma_{1,0} = \gamma_{N,N+1} = 0 \\ f(y_n) = \frac{\delta}{2} \left( \left| y_n + \frac{1}{\delta} \right| - \left| y_n - \frac{1}{\delta} \right| \right).$$

For the computer simulations, we calculate Eq. (1) using the fourth-order Runge-Kutta method with step size  $h = 0.005$ . We set the parameters of this circuit model as follows;  $\alpha_c = 0.460$ ,  $\alpha_p = 0.413$ ,  $\beta = 3.0$  and  $\delta = 470.0$ .

### III. RESULTS

#### A. Fixed the couplings strength

Chaos propagation is researched in ladder chaotic circuits, when the coupling strength is fixed at  $\gamma = 0.002, 0.004, 0.006$ , and  $0.008$ . The horizontal axis is the value of  $x_n$  representing the current flowing in  $L_1$ , and the vertical axis shows the value of  $z_n$  representing the voltage applied to  $C$ .

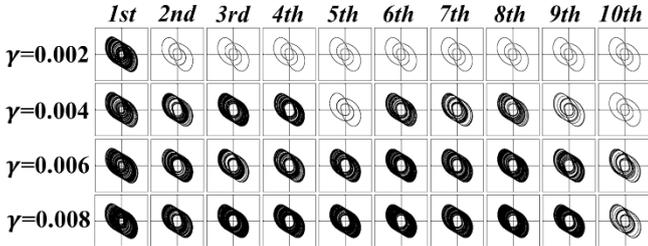


Figure 3: Chaos propagation under fixed coupling strength

The propagation state at almost the same time is shown in Fig. 3. As the coupling strength increases ( $\gamma = 0.004, 0.006$ , and  $0.008$ ), the number of circuits for chaotic propagation becomes larger. In particular, no propagation was observed in case of a coupling strength  $\gamma = 0.002$ .

#### B. Switching coupling strength

The chaos propagation is changed by switching the coupling strength with the difference of  $z_n$  and  $z_{n+1}$  as a threshold. The first method is to switch all the couplings simultaneous, and the second method is to switch each coupling at a certain. How to switch the coupling strength is shown in Fig. 4.

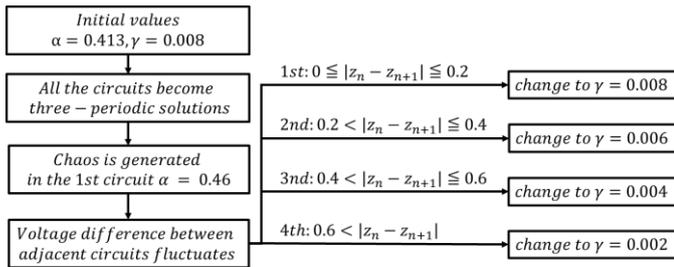


Figure 4: The algorithm for controlling chaotic propagation by coupling strength (The threshold widths are all 0.2).

##### (1) Switching all couplings simultaneous

When the voltage difference at one coupling exceeds the threshold, all couplings switch. In Fig. 5, the maximum number of circuits propagated chaos is fifth, when the third threshold width is 0.1. For the other threshold widths, that is fourth. In a few seconds, the coupling strength decreased from 0.008 to 0.002 and then remained at 0.002.

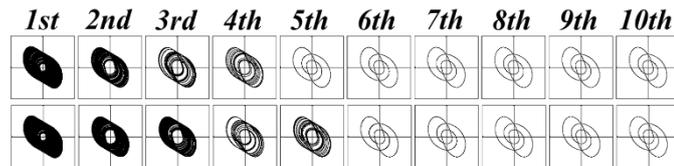


Figure 5: Chaos propagation after a certain time. (Same time as Fig. 3)

##### (2) Switching each coupling

When the coupling strength is switched for each coupling, propagation can be classified into four patterns as shown in the Fig. 6.

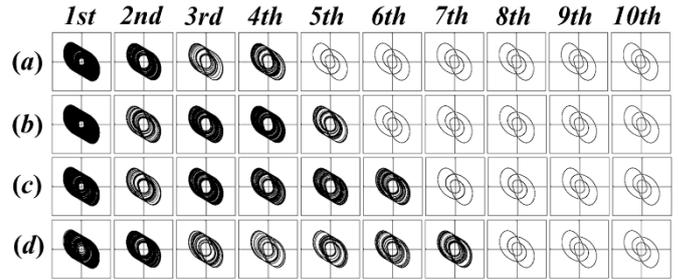


Figure 6: Four patterns of chaos propagation.

When the first threshold width is  $0 \leq |z_n - z_{n+1}| \leq 0.2$ , the chaos propagation can be suppressed as shown in the Fig. 6(a)-(c). As the second threshold width increases to 0.1, 0.2 and 0.3, the number of circuits for chaotic propagation becomes larger to 4th, 5th and 6th as well as when the first threshold widths are 0.1 and 0.3. However, when threshold widths are 0.1, 0.1 and 0.4 from 1st to 3rd step in the algorithm, the number of circuits propagated chaos becomes the largest as shown in Fig. 6(d).

### IV. CONCLUSION

In this study, chaos propagation in a ladder-coupled chaotic circuit was studied when the coupling strength was switched. The propagation was regulated in three ways: fixed coupling strength, simultaneous switching of all the couplings, and switching at each coupling. The results show that the propagation of chaos can be controlled locally, and even in the case of switching at each coupling, the propagation can be suppressed as well as in the method of switching all the couplings at the same time by considering the decision criteria. In the future, we would like to focus on the propagation process such as a recovery from chaos and indirect propagation.

### REFERENCES

- [1] Masahiro WADA, Yoshifumi NISHIO and Akio USHIDA, "Analysis of Bifurcation Phenomena in Two Chaotic Circuits Coupled by an Inductor", IEICE Transactions on Fundamentals, vol. E80-A, no. 5, pp. 869-875, May 1997.
- [2] Yuji TAKAMARU, Yoko UWATE, Thomas OTT and Yoshifumi NISHIO, "Dependence of Clustering Patterns on Density of Chaotic Circuits in Networks", RISP Journal of Signal Processing, vol. 17, no. 4, pp. 103-106, Jul. 2013.
- [3] Yoko UWATE and Yoshifumi NISHIO, "Chaos Propagation in a Ring of Coupled Circuits Generating Chaotic and Three-Periodic Attractors", Proceedings of IEEE Asia Pacific Conference on Circuits and Systems (APCCAS'12), pp. 643-646, Dec. 2012.
- [4] Takahiro CHIKAZAWA, Yoko UWATE and Yoshifumi NISHIO, "Spread of Chaotic Behavior in Scale-Free and Random Networks", Proceedings of Asia Pacific Conference on Postgraduate Research in Microelectronics and Electronics (PrimeAsia'17), pp. 21-24, Oct. 2017.
- [5] Toshiaki NISHIUMI, Yoko UWATE and Yoshifumi NISHIO, "Synchronization Phenomena of Chaotic Circuits with Stochastically-Changed Network Topology", Proceedings of International Symposium on Nonlinear Theory and its Applications (NOLTA'14), pp. 811-814, Sep. 2014.